

# PHYS 705: Classical Mechanics



HW#8 (11/8) and #9 (11/15) move down one more week

HW#8 and #9: CT and Hamilton-Jacobi Eq

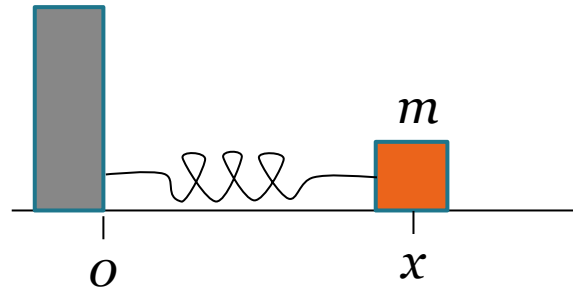
HW#10: Small Oscillations

HW#11: Noninertial Reference Frame and Rigid Body Motion

HW#12: Rigid Body Motion (more practice problems)

## Review from Previous Lecture

## Example: Harmonic Oscillator



$$f(x) = -kx \quad U(x) = \frac{kx^2}{2}$$

$$L = T - U = \frac{m\dot{x}^2}{2} - \frac{kx^2}{2}$$

Since  $U$  does not dep on  $\dot{x}$  and  $x \mapsto x$  does not dep  $t$  explicitly,  $H = E$ .

Define  $\omega = \sqrt{k/m}$  or  $\omega^2 m = k \rightarrow$

$$H = \frac{p^2}{2m} + \frac{m\omega^2}{2} x^2 = \frac{1}{2m} (p^2 + m^2 \omega^2 x^2)$$

Notice that the system is not cyclic with the original variable  $x$ . We will attempt to find a canonical transformation from  $(x, p)$  to  $(X, P)$  such that  $X$  is cyclic so that  $P$  is constant.

## Example: Harmonic Oscillator

With the form of  $H$  as sums of squares  $H = \frac{1}{2m}(p^2 + m^2\omega^2 x^2)$ , we will try to exploit  $\sin^2 + \cos^2 = 1$

Trail Solution:  $p = g(P)\cos(X) \quad x = \frac{g(P)}{m\omega}\sin(X) \quad (*)$

Substituting these into our Hamiltonian, we have

$X$  is cyclic in  $K$  !

$$K = \frac{1}{2m} \left( g^2(P) \cos^2 X + m^2 \omega^2 \frac{g^2(P)}{m^2 \omega^2} \sin^2 X \right) = \frac{g^2(P)}{2m} \quad \leftarrow$$

## Example: Harmonic Oscillator

Need to find the right  $g(P)$  such that the transformation is **canonical**!

Try the following type 1 generating function:  $F = F(x, X, t)$

(with  $x$  and  $X$  being the indep vars)

For  $F$  to be canonical transformation,  $F(x, X, t)$  must satisfy the partial derivative relations (Type 1 in Table 9.1 in G):

$$p = \frac{\partial F}{\partial x}$$

$$P = -\frac{\partial F}{\partial X}$$

## Example: Harmonic Oscillator

Dividing the two trial solutions, we get,

$$\frac{p}{x} = \frac{\cancel{g(P)} \cos X}{\cancel{g(P)} \sin X (1/m\omega)} = m\omega \cot X \quad \text{or} \quad p = m\omega x \cot X$$

Then, the two partial derivative equations give:

$$p = \frac{\partial F}{\partial x} = m\omega x \cot X \quad \longrightarrow \quad F = F(x, X, t) = \frac{m\omega x^2}{2} \cot X$$

$$P = -\frac{\partial F}{\partial X} \quad \longrightarrow \quad P = \frac{m\omega x^2}{2} \frac{1}{\sin^2 X}$$

⋮

## Example: Harmonic Oscillator

•  
•  
•

$$x = \sqrt{\frac{2P}{m\omega}} \sin X \qquad p = \sqrt{2m\omega P} \cos X$$

By comparing with our trial transformation equation,

$$\begin{cases} x = \frac{g(P)}{m\omega} \sin X \\ p = g(P) \cos X \end{cases} \quad \text{gives} \quad g(P) = \sqrt{2m\omega P} \quad (*)'$$

Then, the transformed Hamiltonian is:  $K = \frac{g^2(P)}{2m} = \frac{2m\omega P}{2m} = \omega P$



## Example: Harmonic Oscillator

$$\frac{\partial F}{\partial t} = 0$$

With the transformed Hamiltonian  $K = \omega P = H = E$ , the Hamilton's Equations give,

$$\dot{P} = -\frac{\partial K}{\partial X} = 0 \quad (X \text{ is } \text{cyclic})$$

$$\rightarrow \boxed{P = \text{const}}$$

$$\dot{X} = \frac{\partial K}{\partial P} = \omega$$

$$\rightarrow \boxed{X = \omega t + \alpha}$$

depends on IC

Together with the inverse transform,  $x(t)$ ,  $p(t)$  can be solved explicitly,

$$x = \sqrt{\frac{2P}{m\omega}} \sin(\omega t + \alpha)$$

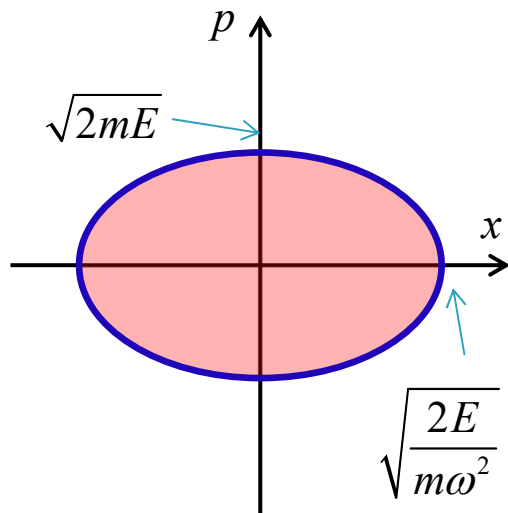
$$p = \sqrt{2m\omega P} \cos(\omega t + \alpha)$$

## Example: Harmonic Oscillator

Then with  $K = \omega P = E$

$$x = \sqrt{\frac{2E}{m\omega^2}} \sin(\omega t + \alpha)$$

$$p = \sqrt{2mE} \cos(\omega t + \alpha)$$

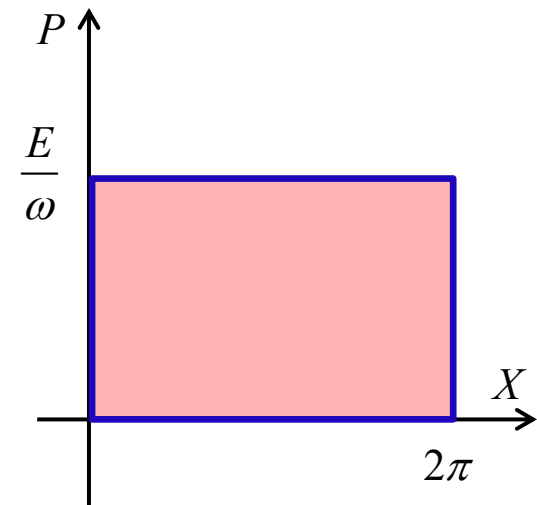


Phase Space area  
is invariant!

$$area = \frac{2\pi E}{\omega}$$

$(X, P)$  is an example of  
Action-Angle variable pair,

$$X = \omega t + \alpha \quad P = const$$



## “Symplectic” Approach & Poisson Bracket

By considering  $Q_j(q, p)$  and  $P_j(q, p)$  as function of  $q$  and  $p$  explicitly and they satisfy the Hamilton's Equations,

We can derive the following four “direct conditions” for various  $(i, j)$  for the canonical transformation pair,  $Q_j(q, p), P_j(q, p)$  .

$$\begin{aligned} \left( \frac{\partial Q_j}{\partial q_i} \right)_{q,p} &= \left( \frac{\partial p_i}{\partial P_j} \right)_{Q,P} & \left( \frac{\partial Q_j}{\partial p_i} \right)_{q,p} &= - \left( \frac{\partial q_i}{\partial P_j} \right)_{Q,P} \\ \left( \frac{\partial P_j}{\partial q_i} \right)_{q,p} &= - \left( \frac{\partial p_i}{\partial Q_j} \right)_{Q,P} & \left( \frac{\partial P_j}{\partial p_i} \right)_{q,p} &= \left( \frac{\partial q_i}{\partial Q_j} \right)_{Q,P} \end{aligned}$$